

Egzotikus elektromágneses jelenségek alacsony hőmérsékleten

Mihály György
BME Fizikai Intézet

Hall effektus

Edwin Hall és az „összenyomhatatlan elektromosság”
Kvantum Hall effektus

Mágneses áram

anomális Hall effektus
spin forgatónyomaték

Spin-polarizáció mérése

Andrejev reflexió
Mágneses mikroszkópia nanométeres tartományban

<http://dept.phy.bme.hu/egzotikus.pdf>

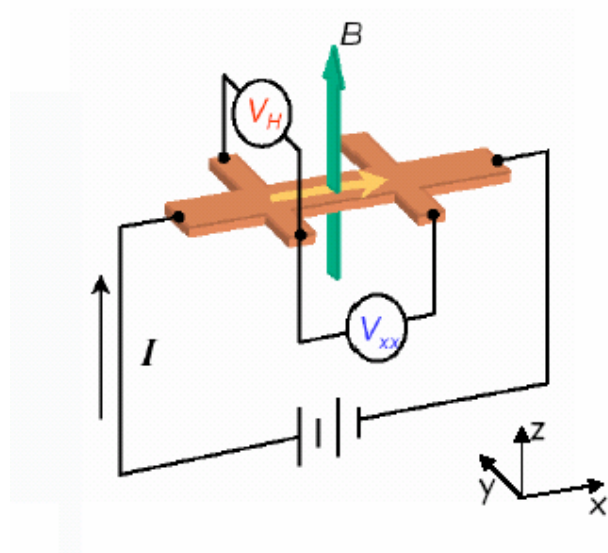
Hall effektus

Hall jelenség



Edwin Hall (1855-1938)

American Journal of Mathematics, vol 2 (1879)



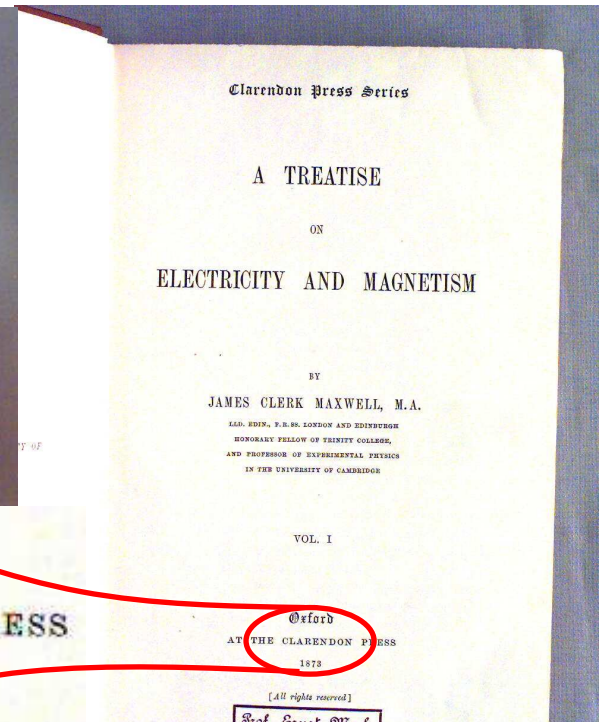
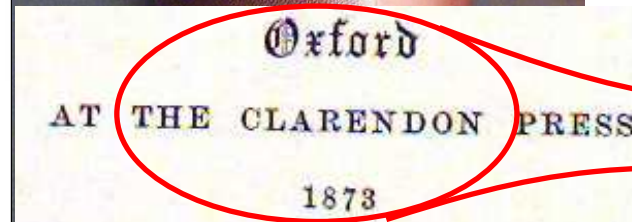
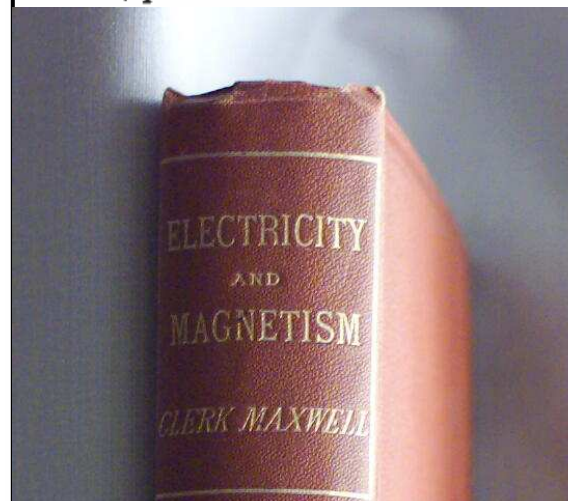
Hall-állandó

$$R_H = \frac{E_y}{j_x B_z} = \frac{1}{ne}$$

On a New Action of the Magnet on Electric Currents.

By E. H. HALL, *Fellow of the Johns Hopkins University.*

SOMETIME during the last University year, while I was reading Maxwell's Electricity and Magnetism in connection with Professor Rowland's lectures, my attention was particularly attracted by the following passage in Vol. II, p. 144:



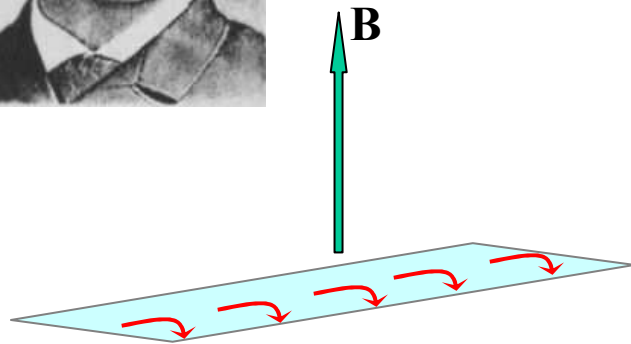
Hall effektus

Hall jelenség: Lorentz erő



Edwin Hall (1855-1938)

American Journal of Mathematics, vol 2 (1879)



Lorentz erő:

$$\vec{F} = -e[\vec{E} + \vec{v} \times \vec{B}]$$

On a New Action of the Magnet on Electric Currents.

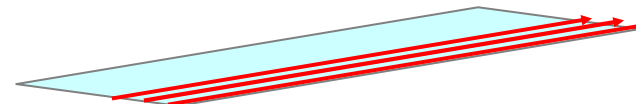
By E. H. HALL, *Fellow of the Johns Hopkins University.*

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"It must be carefully remembered, that the mechanical force which urges a conductor carrying a current across the lines of magnetic force, acts, not on the electric current, but on the conductor which carries it. If the

Finding these two authorities at variance, I brought the question to Prof. Rowland. He told me he doubted the truth of Maxwell's statement and had sometime before made a hasty experiment for the purpose of detecting, if possible, some action of the magnet on the current itself, though without success. Being very busy with other matters however, he had no immediate intention of carrying the investigation further.

Prof. Rowland hipotézise



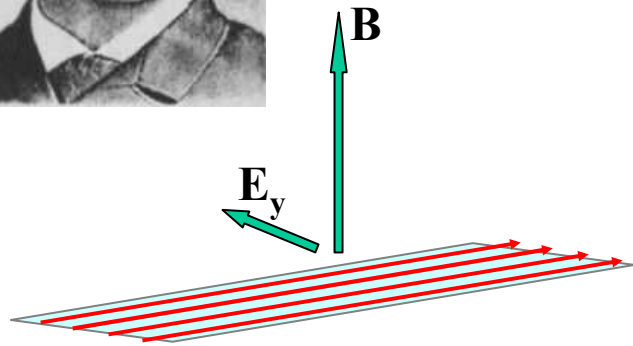
Hall effektus

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But though conclusive, apparently, in respect to any change of resistance, the above experiments are not sufficient to prove that a magnet cannot affect an electric current. If electricity is assumed to be an incompressible fluid, as some suspect it to be, we may conceive that the current of electricity flowing in a wire cannot be forced into one side of the wire or made to flow in any but a symmetrical manner. The magnet may tend to deflect the current without being able to do so. It is evident, however, that in this case there would exist a state of stress in the conductor, the electricity pressing, as it were, toward one side of the wire. Reasoning thus, I thought it necessary, in order to make a thorough investigation of the matter, to test for a difference of potential between points on opposite sides of the conductor.

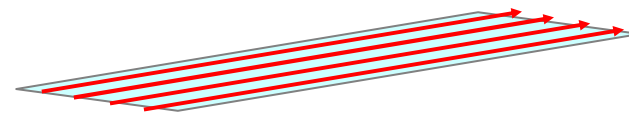
Lorentz erő:

$$\vec{F} = -e[\vec{E} + \vec{v} \times \vec{B}] \quad \vec{j} = -ne\vec{v}$$

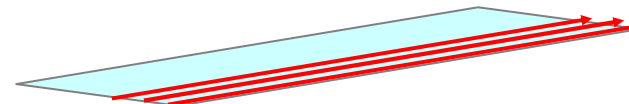
Hall-állandó

$$R_H = \frac{E_y}{j_x B_z} = \frac{1}{ne}$$

Edwin Hall elképzelése



Prof. Rowland hipotézise



Összenyomhatatlan „elektromosság” - kvantummechanika

Schrödinger egyenlet

$$-\frac{\hbar^2}{2m} \nabla^2 \psi = \varepsilon \psi$$

síkhullám

$$\psi(\mathbf{r}) = \frac{1}{\sqrt{V}} e^{i\mathbf{k}\mathbf{r}}$$

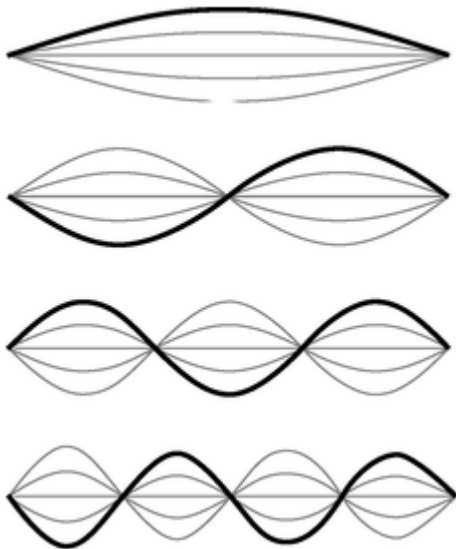
energia

$$\varepsilon(\mathbf{k}) = \frac{\hbar^2 \mathbf{k}^2}{2m}$$

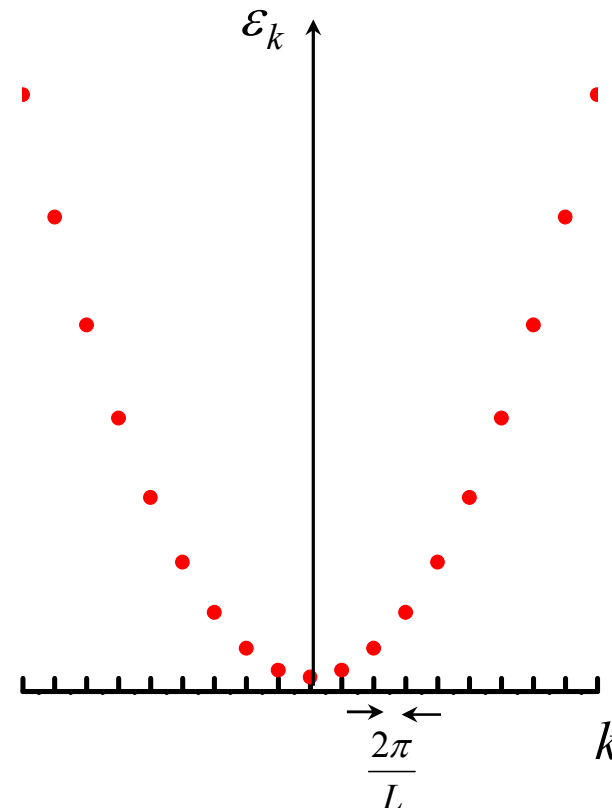
impulzus

$$\mathbf{p} = \hbar \mathbf{k}$$

A hullámszám lehetséges értékei



$$k = \frac{2\pi}{L} n$$



Összenyomhatatlan „elektromosság” - kvantummechanika

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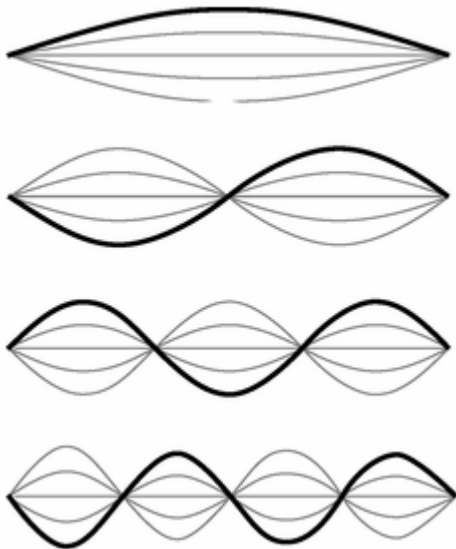
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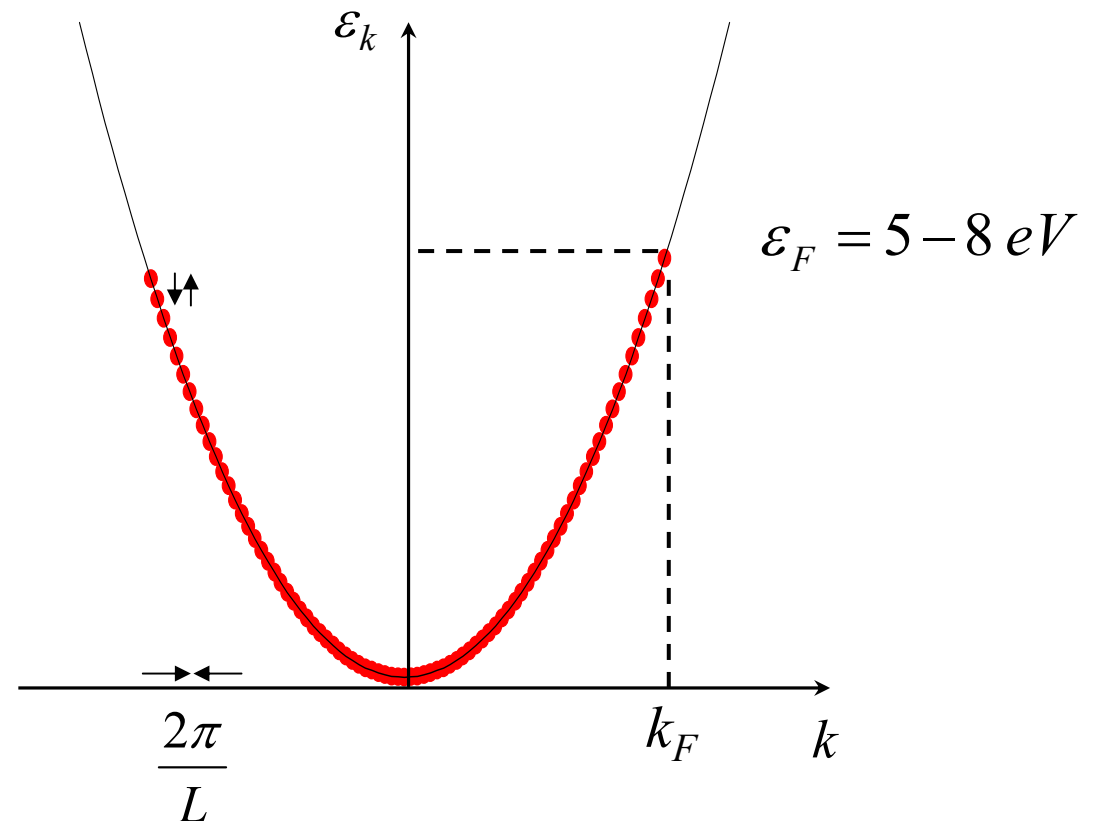
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Sok elektron - Pauli elv



Alapállapot



Összenyomhatatlan „elektromosság” - kvantummechanika

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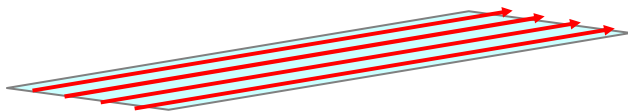
$$\varepsilon(\mathbf{k}) = \frac{\hbar^2 \mathbf{k}^2}{2m}$$

impulzus

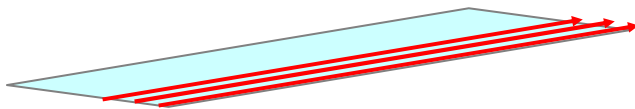
$$\mathbf{p} = \hbar \mathbf{k}$$

Ugyanannyi elektron kisebb helyen:
sokkal nagyobb kinetikus energia

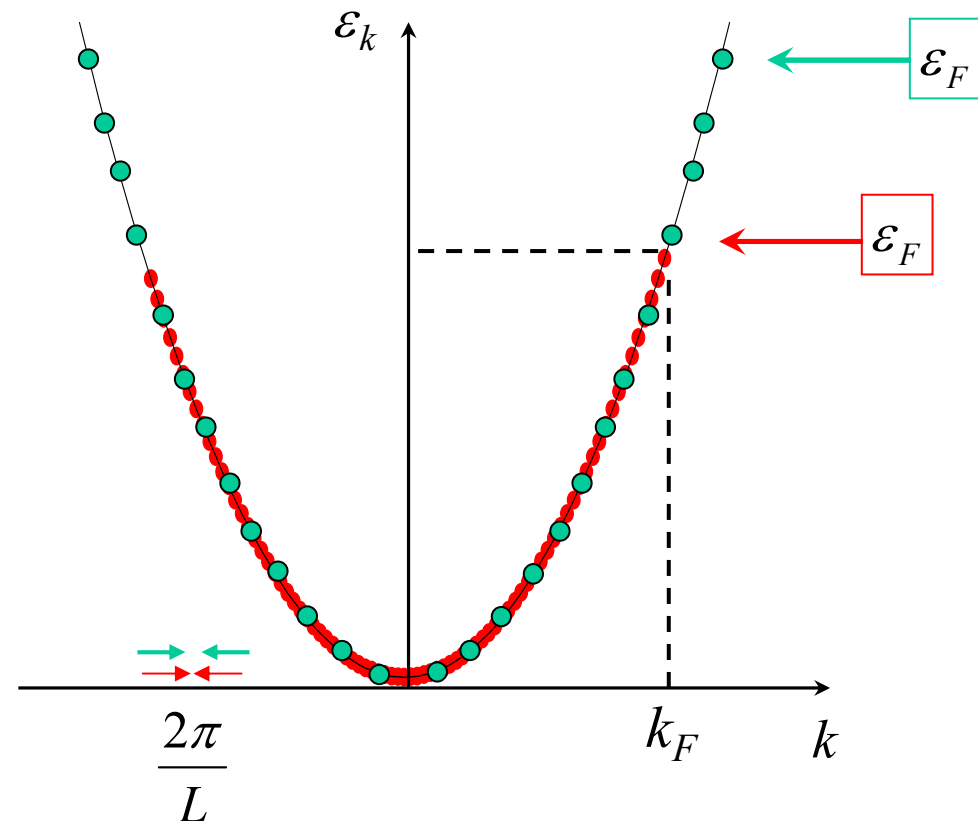
Edwin Hall elképzelése



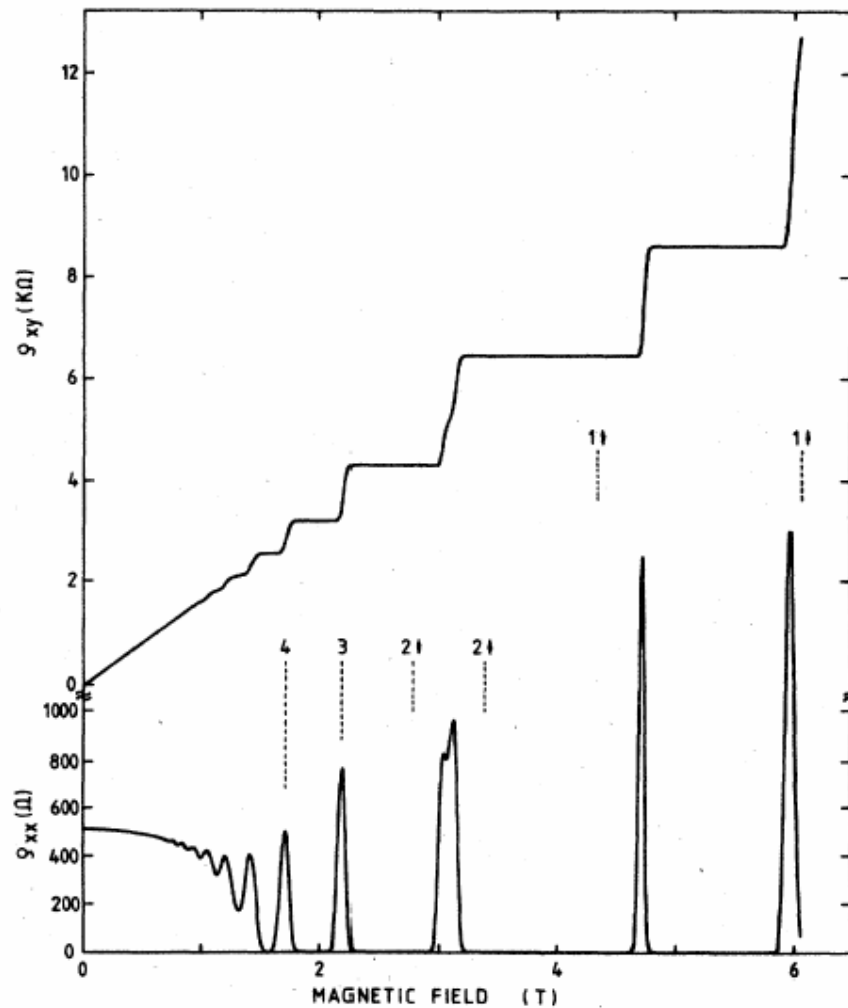
Prof. Rowland hipotézise



Alapállapot



Kvantum Hall effektus



Klaus von Klitzing
Nobel-díj, 1985

FIG. 14. Experimental curves for the Hall resistance $R_H = \rho_{xy}$ and the resistivity $\rho_{xx} \sim R_x$ of a heterostructure as a function of the magnetic field at a fixed carrier density corresponding to a gate voltage $V_g = 0$ V. The temperature is about 8 mK.

T=8 mK

Kvantum Hall effektus

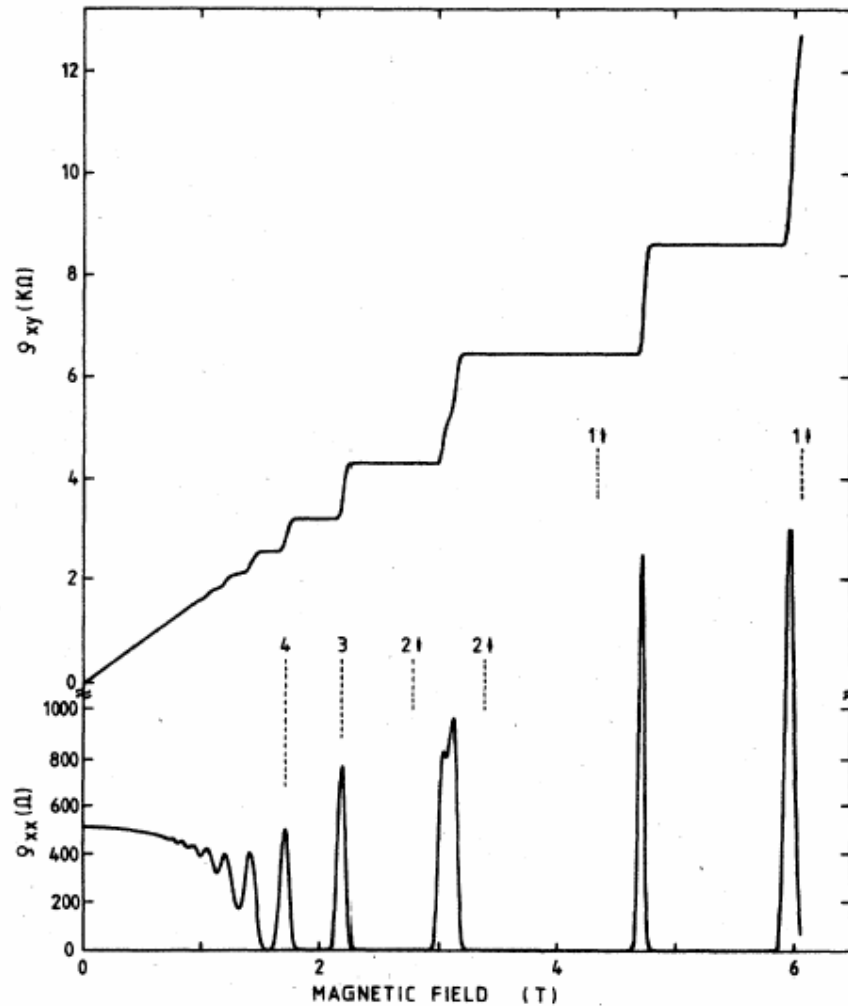
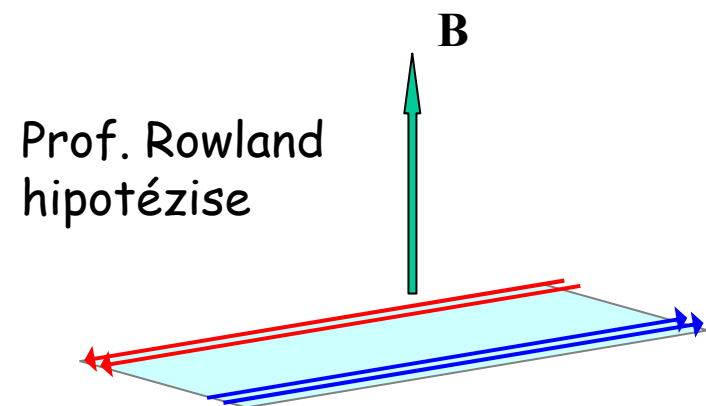


FIG. 14. Experimental curves for the Hall resistance $R_H = \rho_{xy}$ and the resistivity $\rho_{xx} \sim R_x$ of a heterostructure as a function of the magnetic field at a fixed carrier density corresponding to a gate voltage $V_g = 0$ V. The temperature is about 8 mK.

K. von Klitzing, Rev. Mod. Phys **58**, 519 (1986)



Henry Augustus Rowland
(1848-1901)



Kvantum Hall effektus

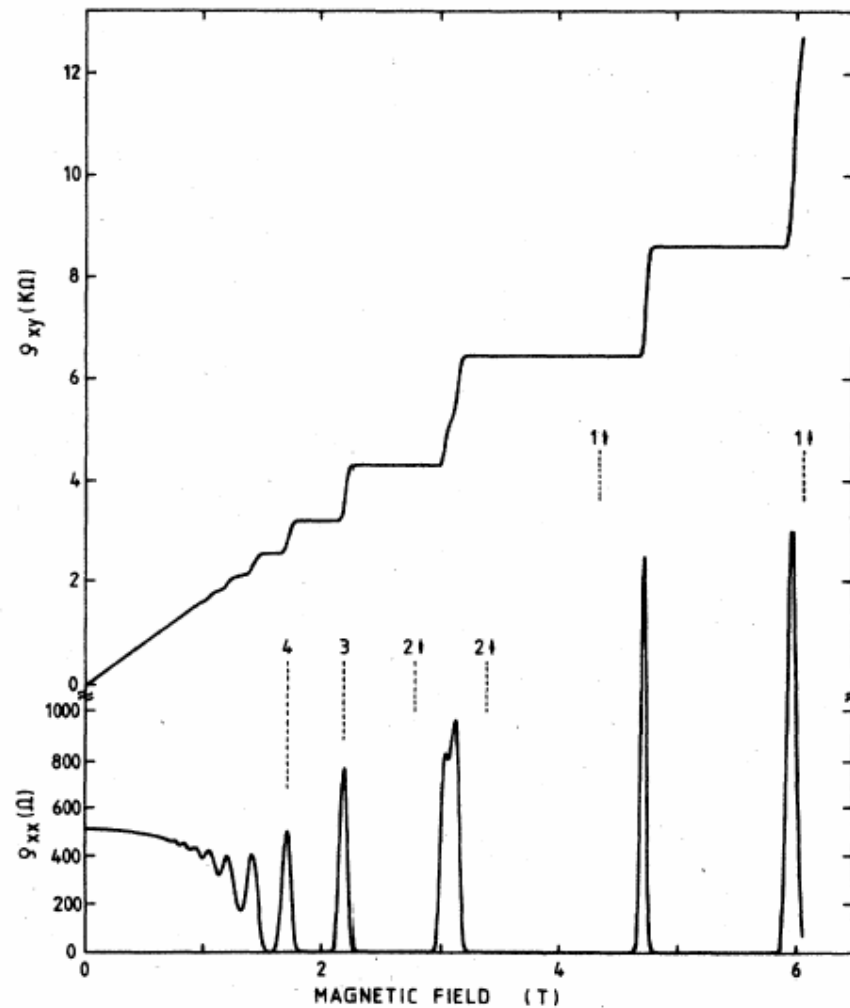
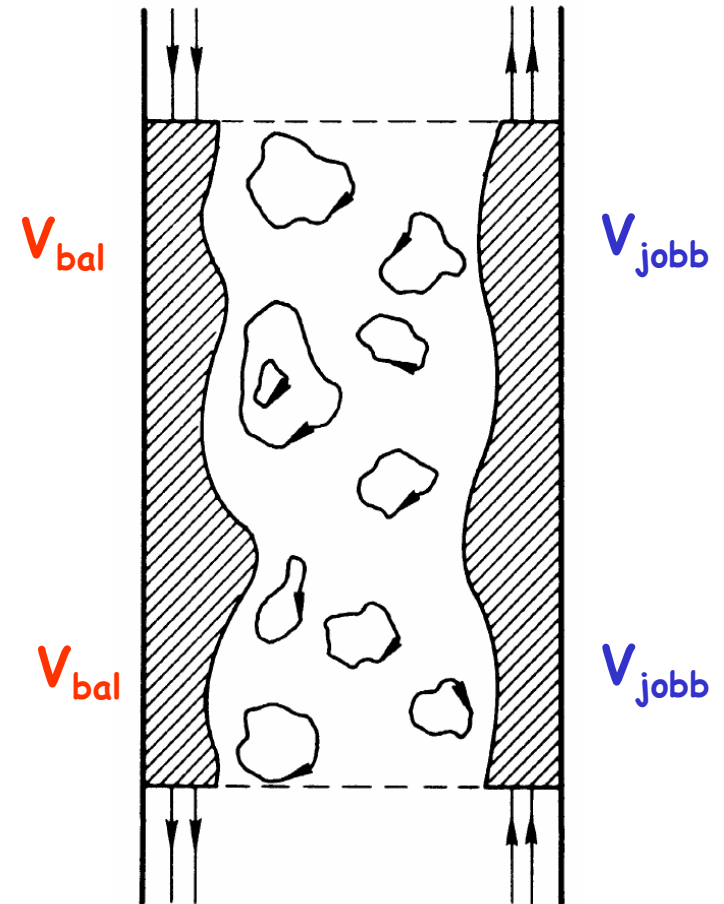
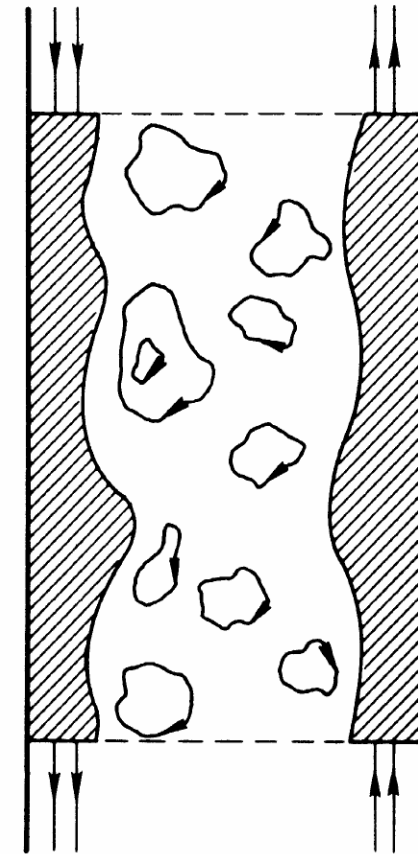
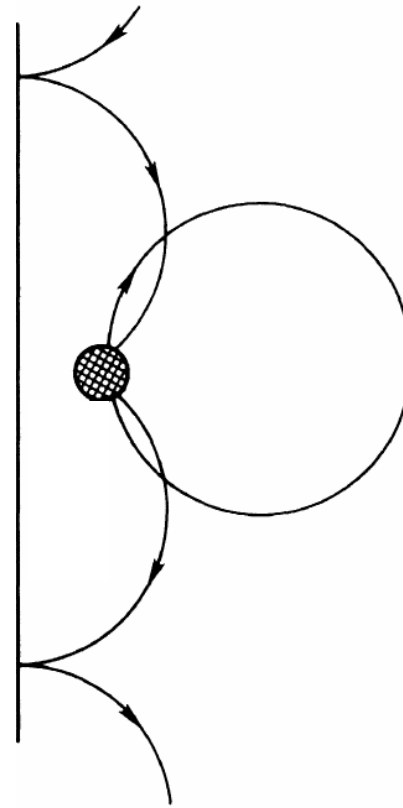
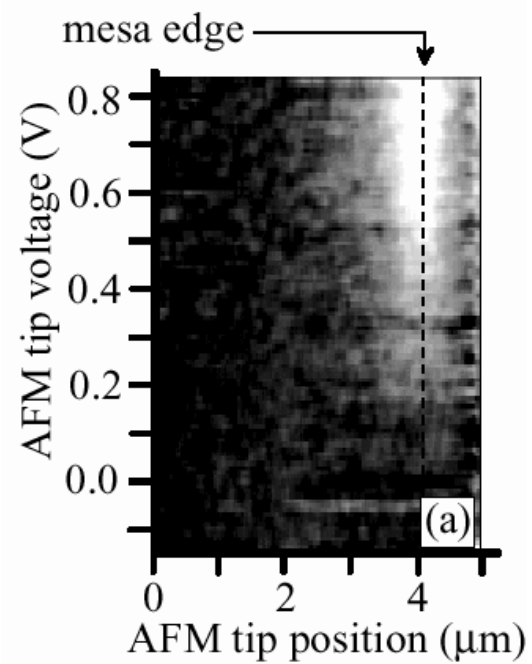


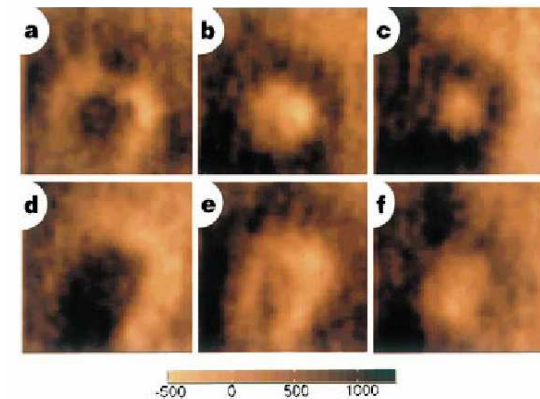
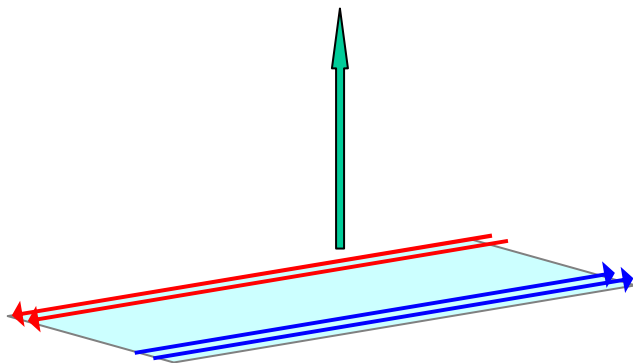
FIG. 14. Experimental curves for the Hall resistance $R_H = \rho_{xy}$ and the resistivity $\rho_{xx} \sim R_x$ of a heterostructure as a function of the magnetic field at a fixed carrier density corresponding to a gate voltage $V_g = 0$ V. The temperature is about 8 mK.



Az áramot ellentétes irányban
szállító elektronok
nem találkoznak



N. B. Zhitenev et al. Nature **404**, 473 (2000)



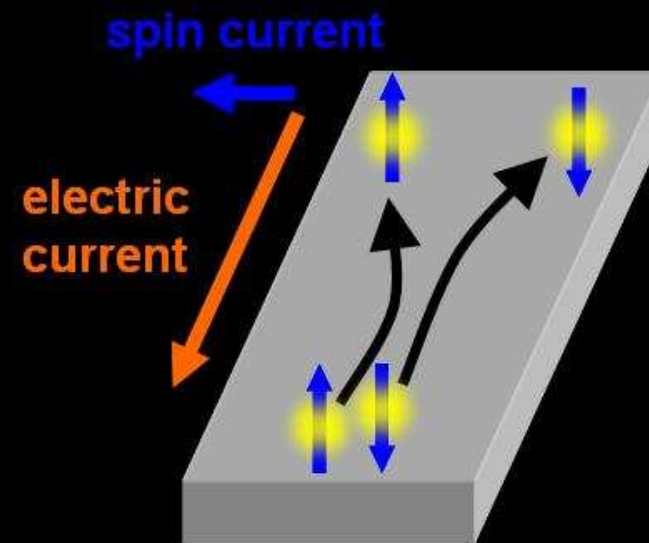
Paul Dirac
1928
relativisztikus kvantummechanika



Elektron spin

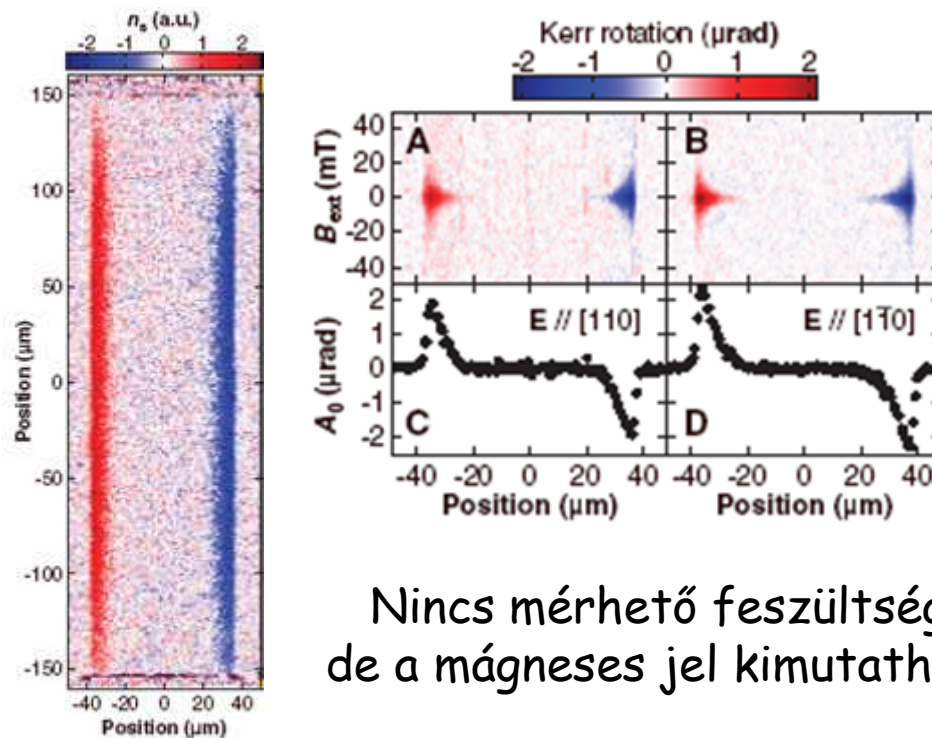
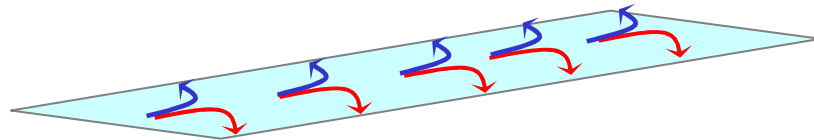


Nobel-díj 1933



spin Hall effektus kísérleti kimutatása, 2004

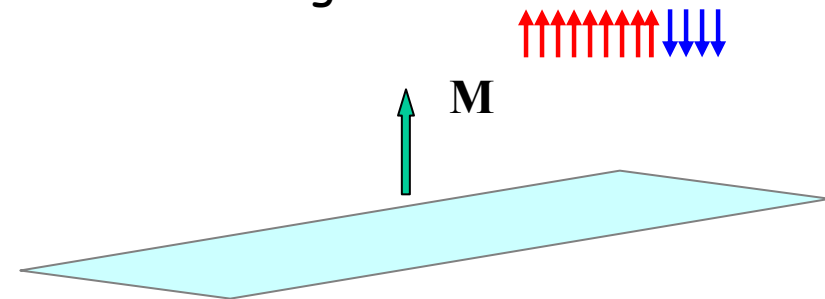
spin Hall effektus
spin állapot szerint kétféle elektron, $\uparrow\downarrow$:
eltérő terjedés a kristályban



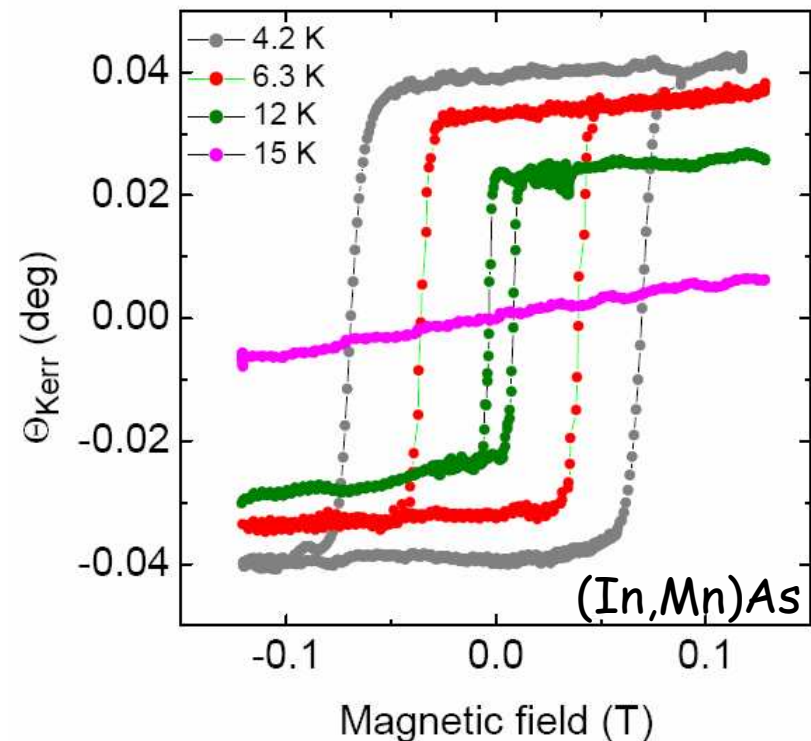
Nincs mérhető feszültség,
de a mágneses jel kimutatható

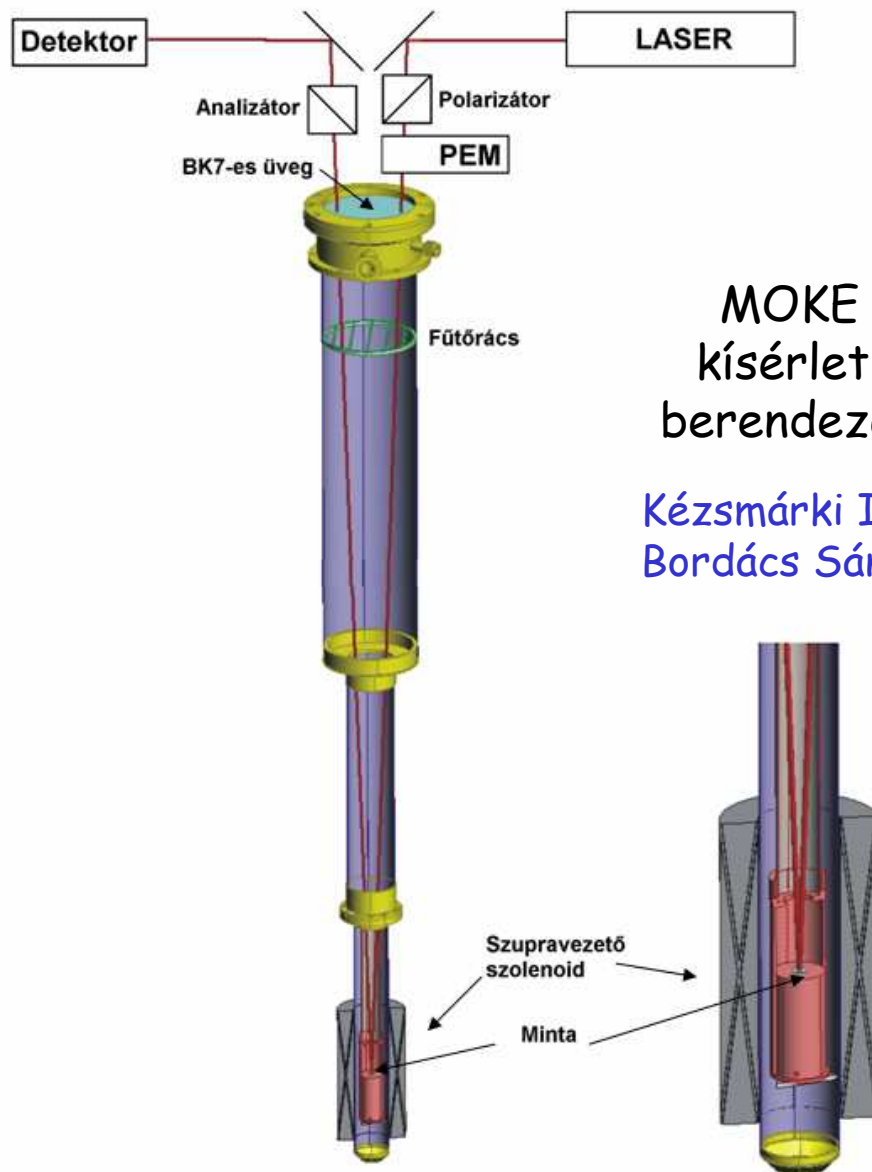
Y.K. Kato *et al.*, Scienceexpress 11 Nov.(2004)

Mágneses minta

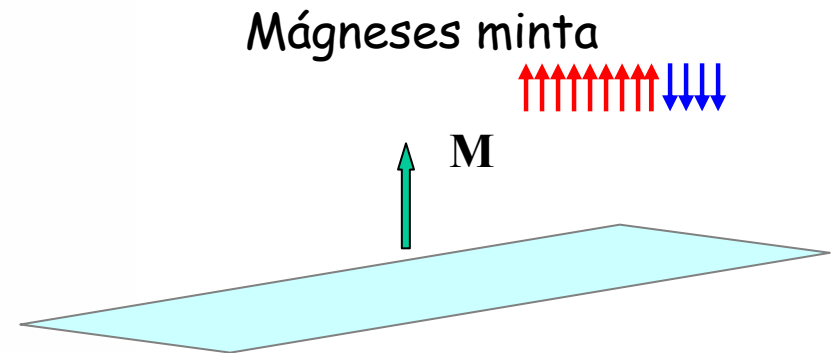


A mágneszettség meghatározása
optikai Kerr-forgatás mérésével

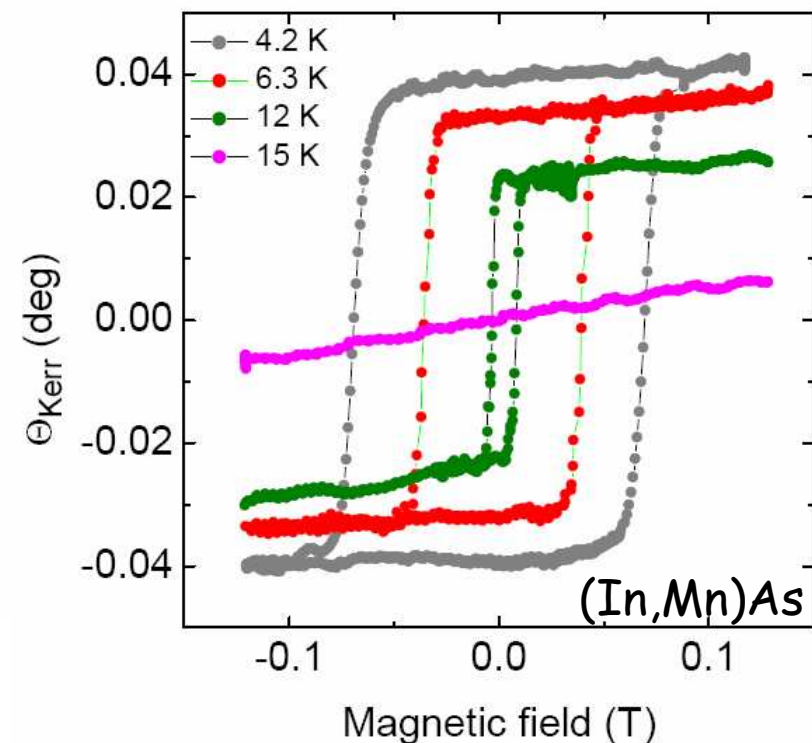




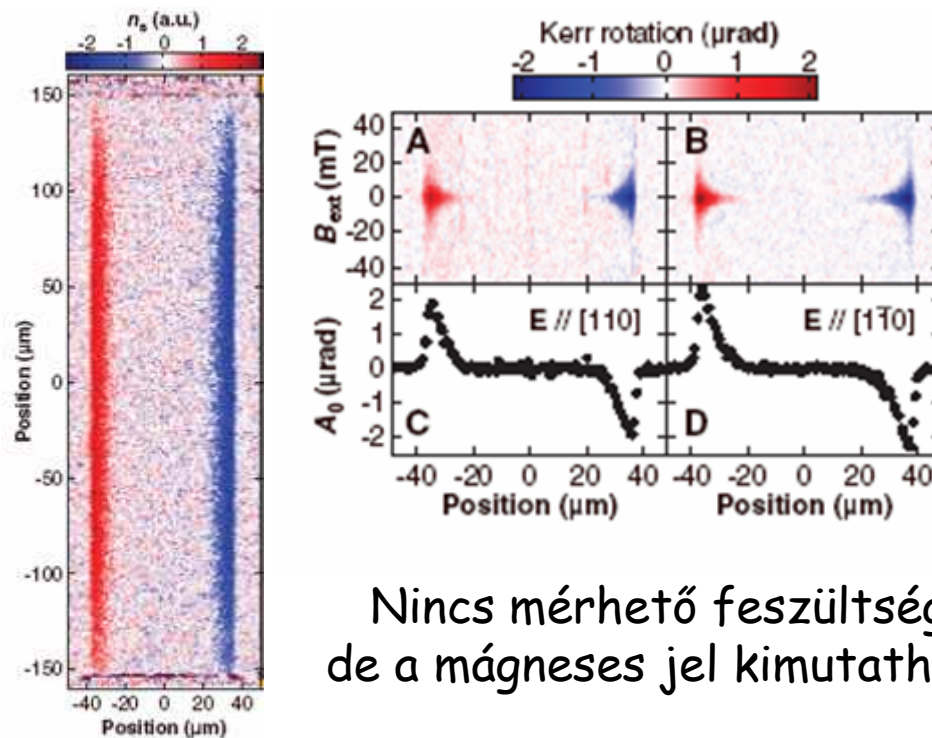
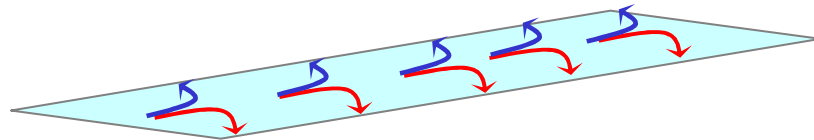
Szupravezető mágnes $B_{\max} = 14 \text{ T}$
Mágneszettség felbontás $\sim 5 \times 10^{-9} \text{ emu}$



A mágneszettség meghatározása
optikai Kerr-forgatás mérésével



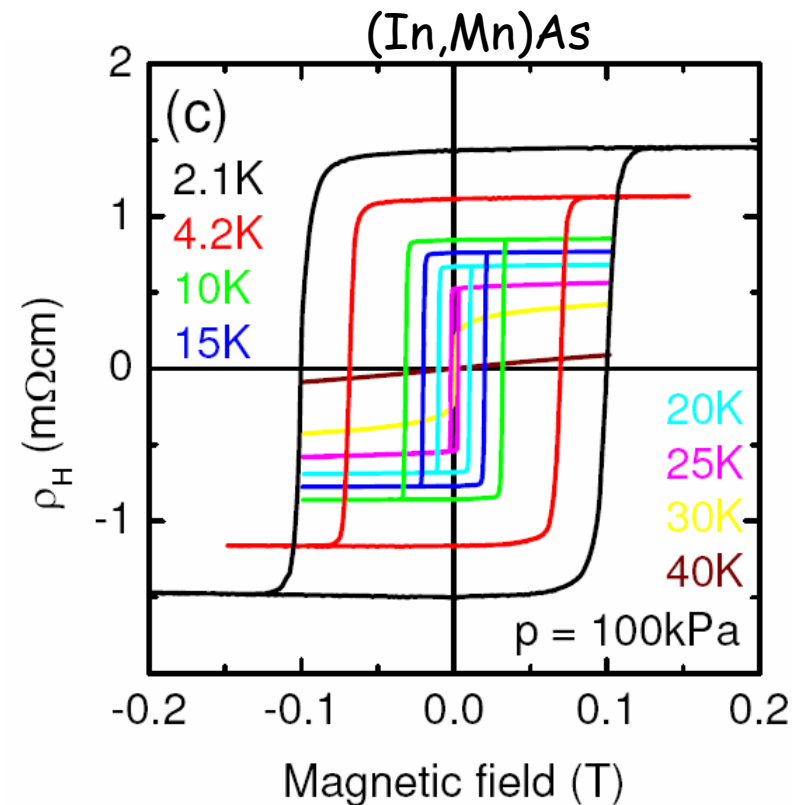
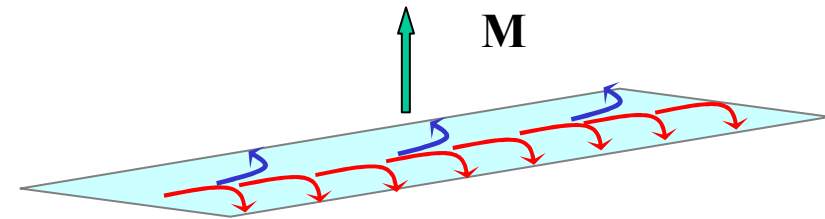
spin Hall effektus
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Nincs mérhető feszültség,
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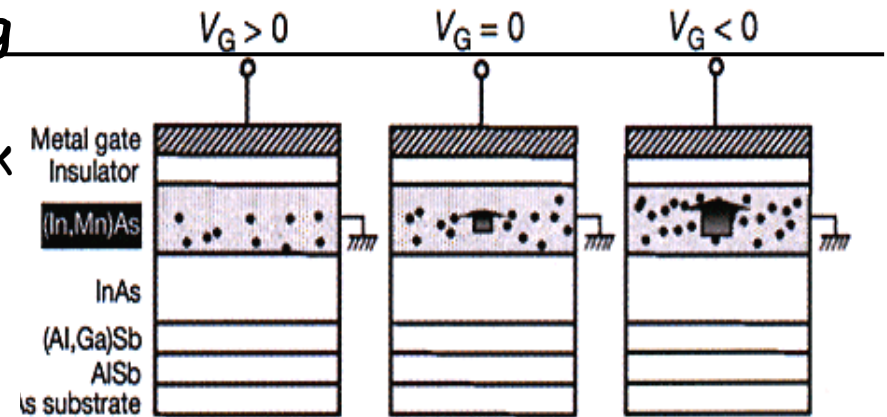
Y.K. Kato *et al.*, *Scienceexpress* 11 Nov.(2004)

Mágneses minta
nem egyforma spin-betöltés:
keresztirányú feszültség mérhető



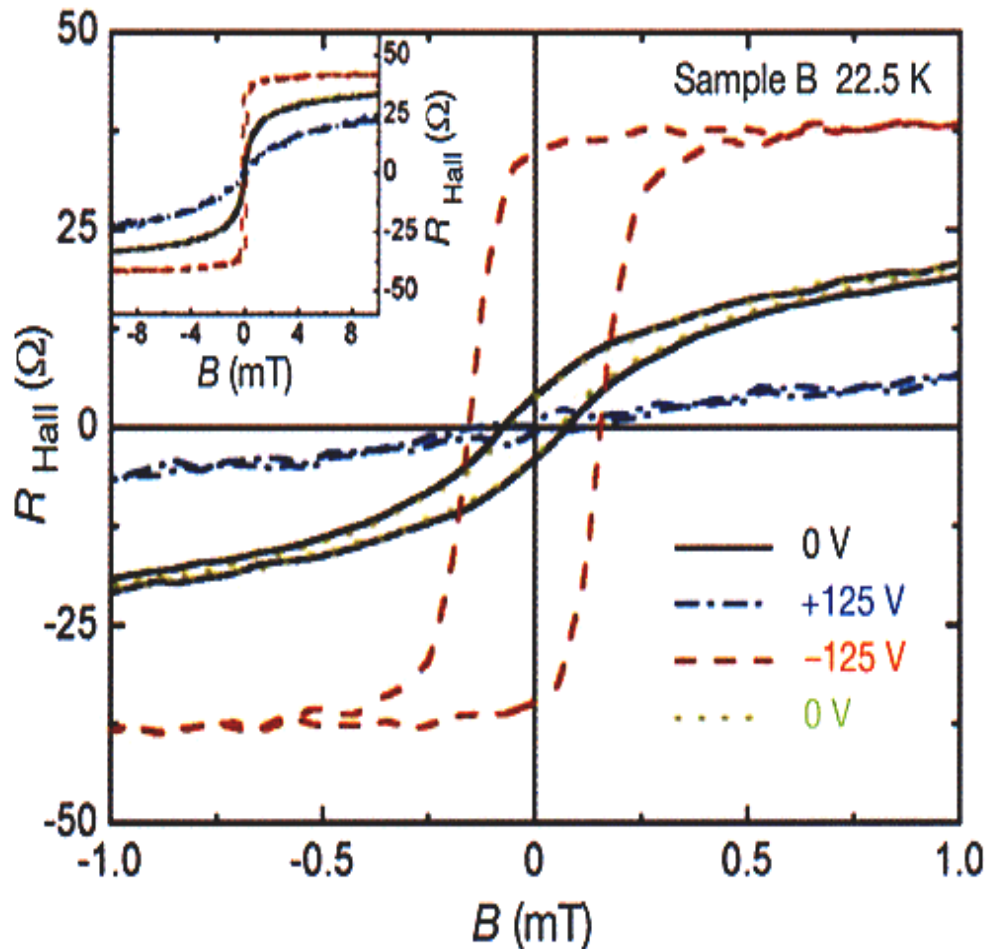
Elektromos térrel kontrollált mágnesség

1. a mágneses csatolást a lyukak közvetítik
2. a lyukak száma FET szerkezetben változtatható

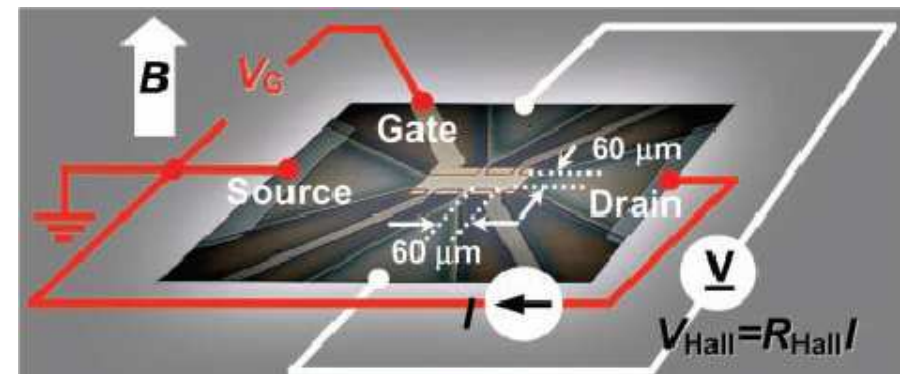


Elektromosan vezérelt mágneses memória

Kontroll -- kapufeszültség
Detektálás -- anomális Hall



H. Ohno *et al.*, Nature **408**, 944(2000)

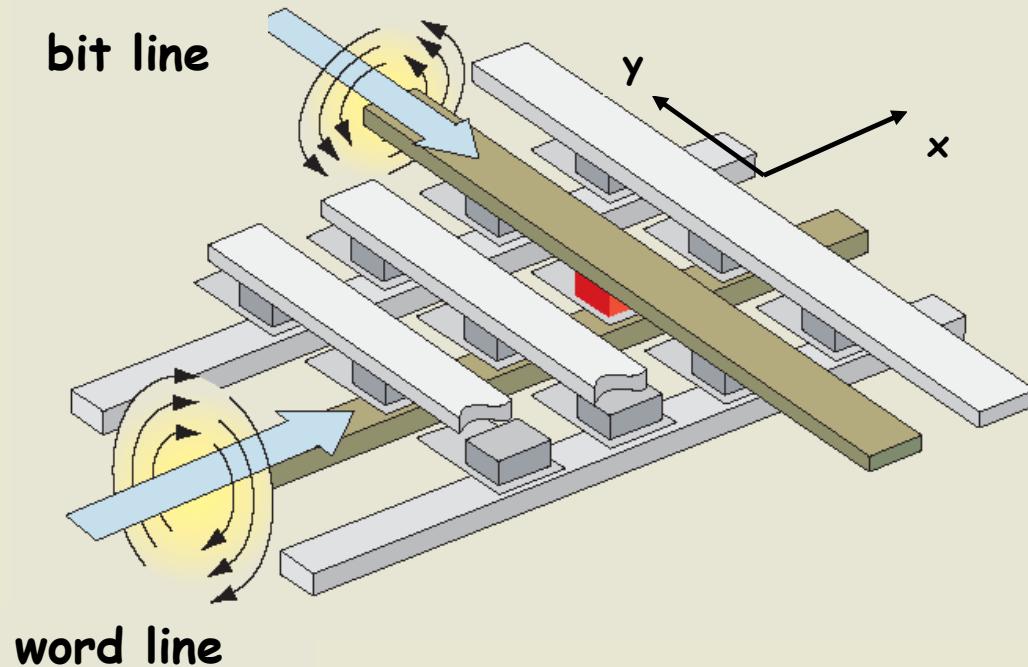


D. Chiba *et al.*, Science **301**, 943 (2003)

Spintronika - Nobel díj 2007

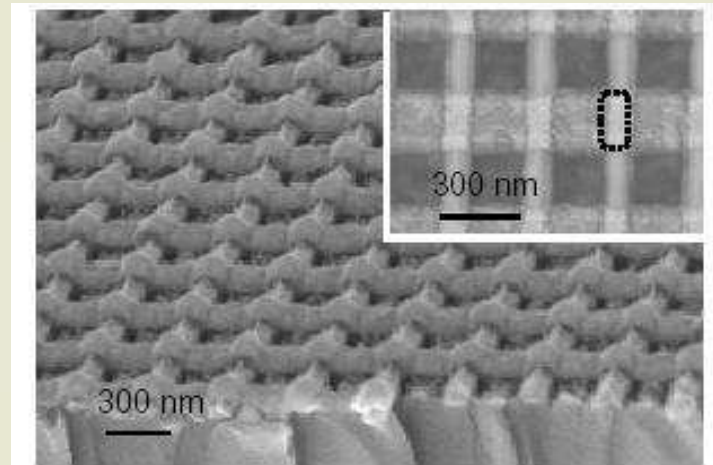
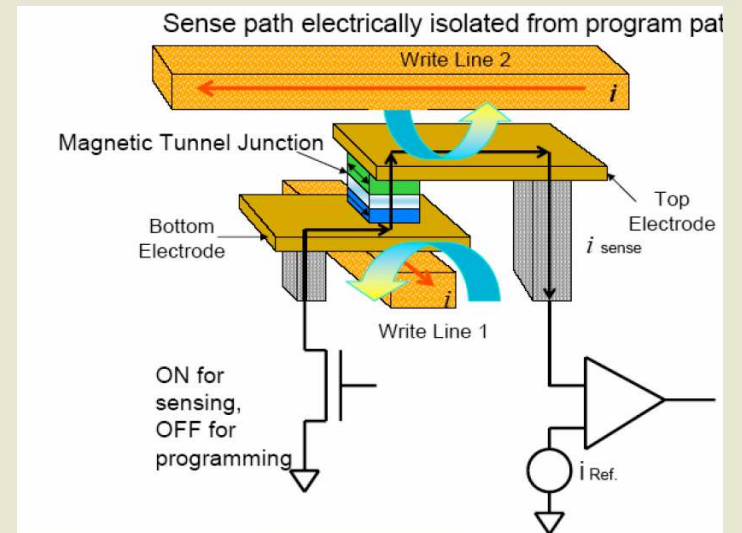
A.Fert, P. Grünberg

MRAM (Magnetic Random Access Memory)

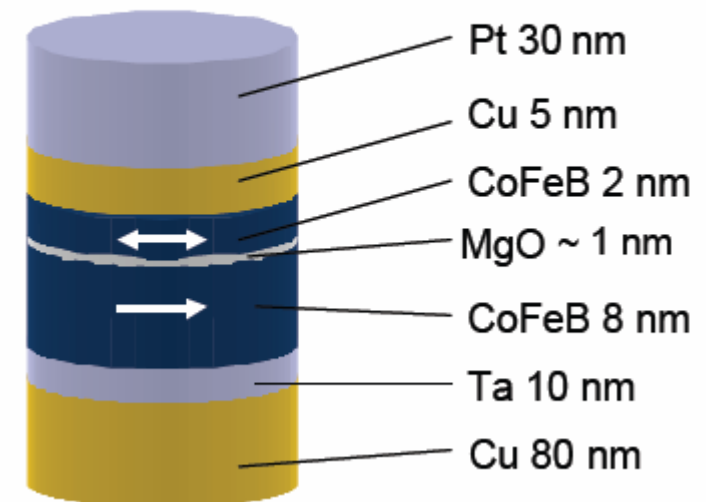
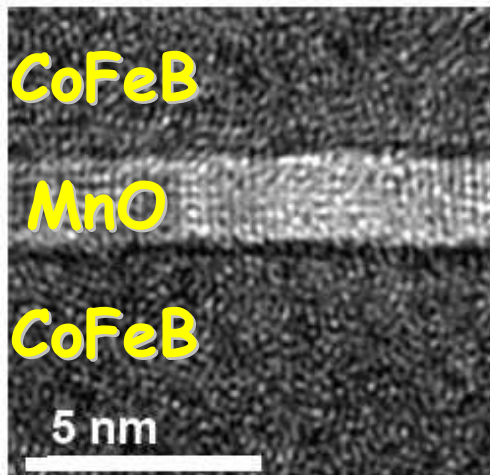
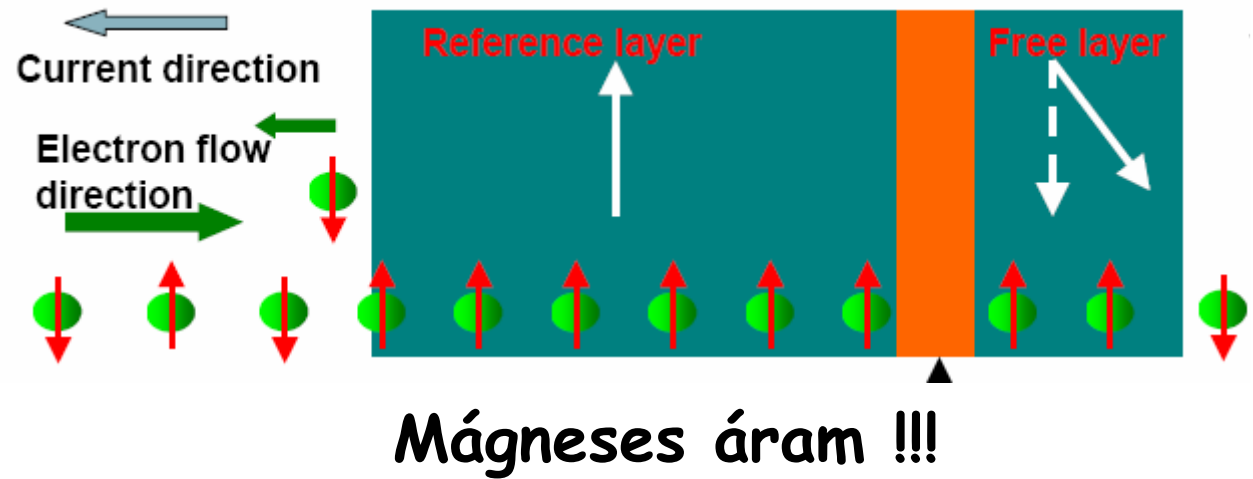
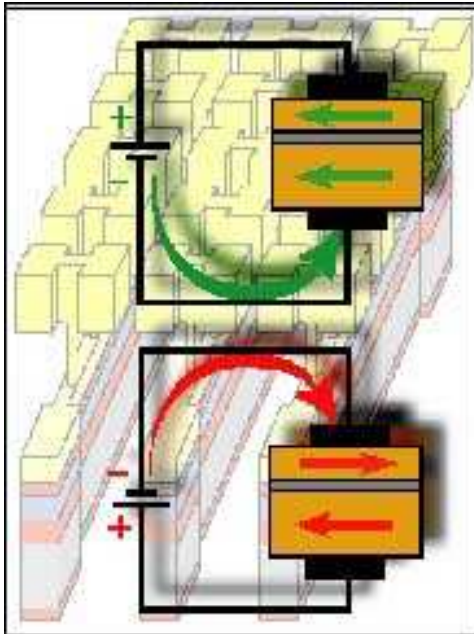


A „bit line” áramszintje alacsonyabb, mint ami szabad réteg mágnesezettségének átfordításához szükséges kivéve, ha a „word line”-ra kapcsolt áram tere

non-volatile
1000-szer gyorsabb írás,
mint a flash memory-ban
adatok elérése 1/10000
része, mint a HDD-ben



Spin fogatónyomaték



Yousa *et al.*, Nature Materials 3, 868 (2004)

Toshiba develops new MRAM device which opens the way to giga-bits capacity

06 November, 2007

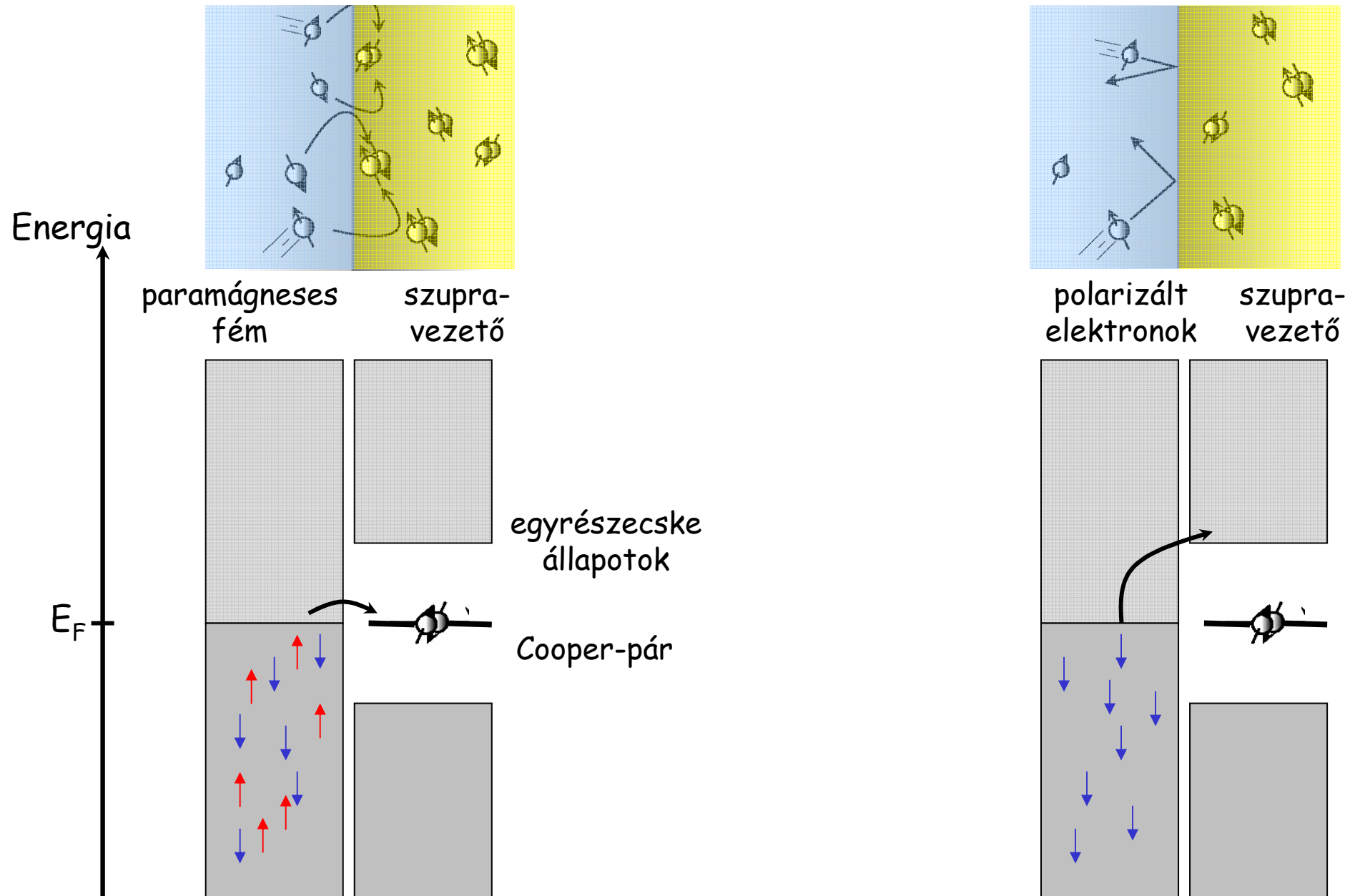
World's first perpendicular magnetic anisotropy MTJ device

Tokyo--Toshiba Corporation today announced important breakthroughs in key technologies for magnetoresistive random access memory (MRAM), a promising, next-generation semiconductor memory device. The company has successfully fabricated a MRAM memory cell integrating the new technologies and verified its stable performance.

Spin transfer switching uses the properties of electron spin to invert magnetization and writes data at very low power levels. It is widely regarded as a major candidate among next-generation principles for new memory devices.

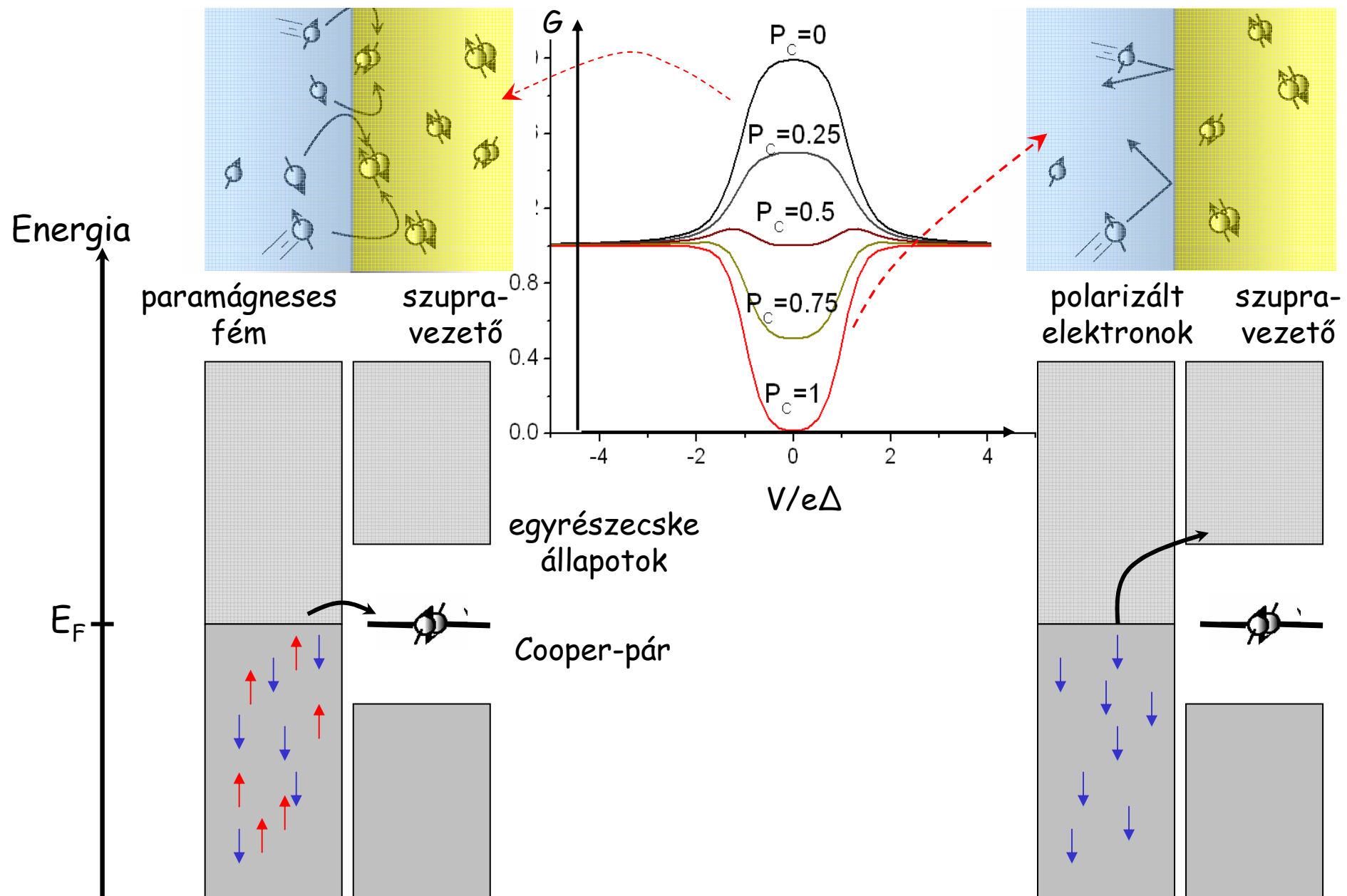
Andrejev spektroszkópia – az elektronok spin polarizációjának mérése

Normal elektron $\rightarrow$ Cooper pár konverzió



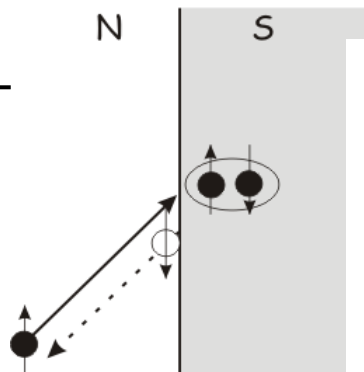
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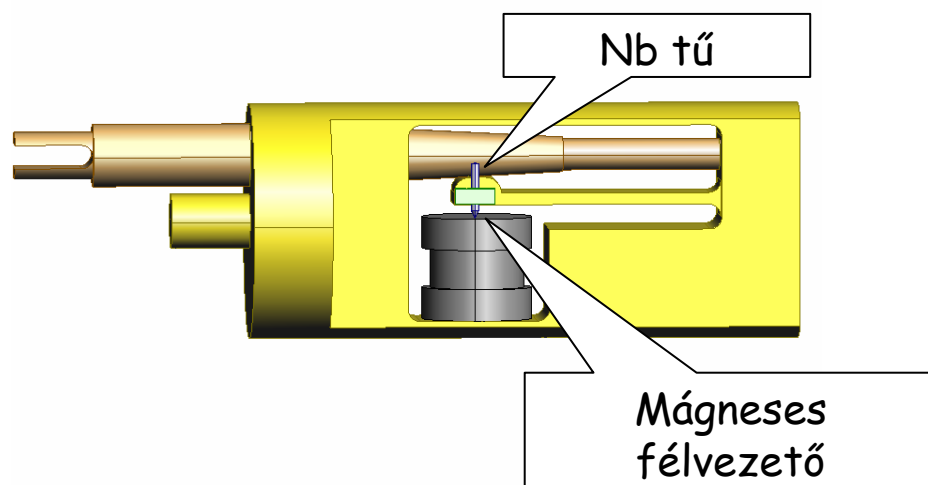


Andrejev reflexió

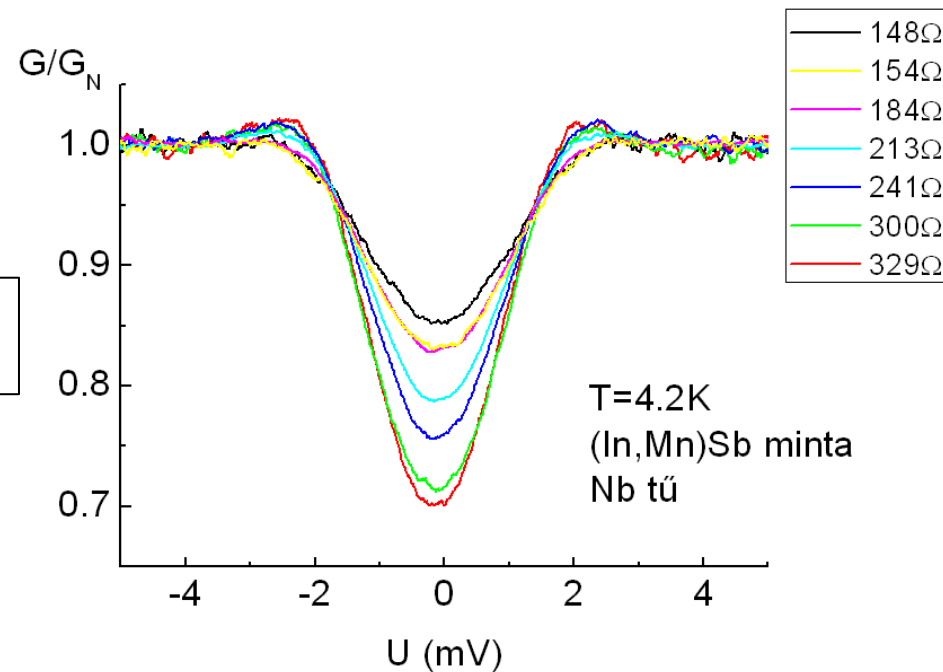
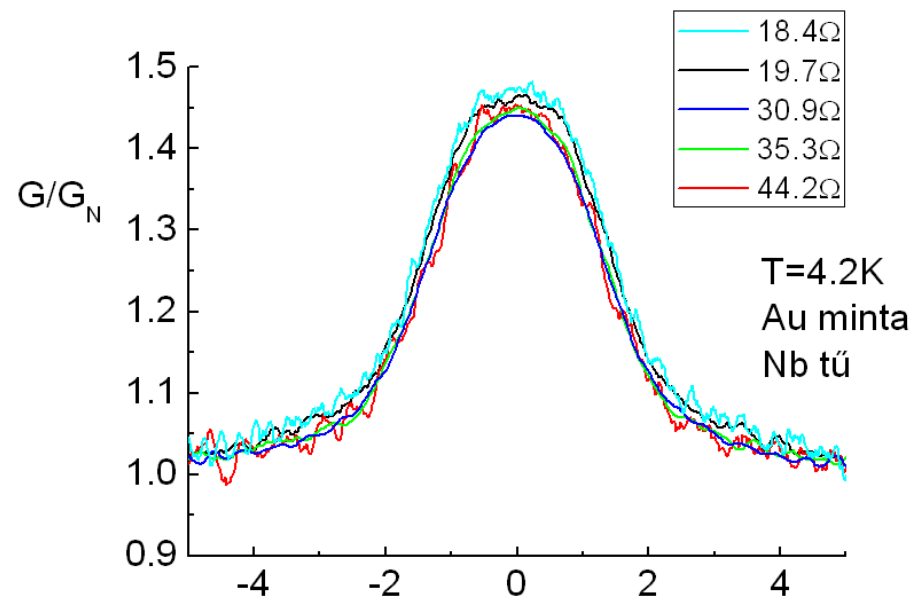
A.F. Andrejev:
JETP 19 1228 (1964)



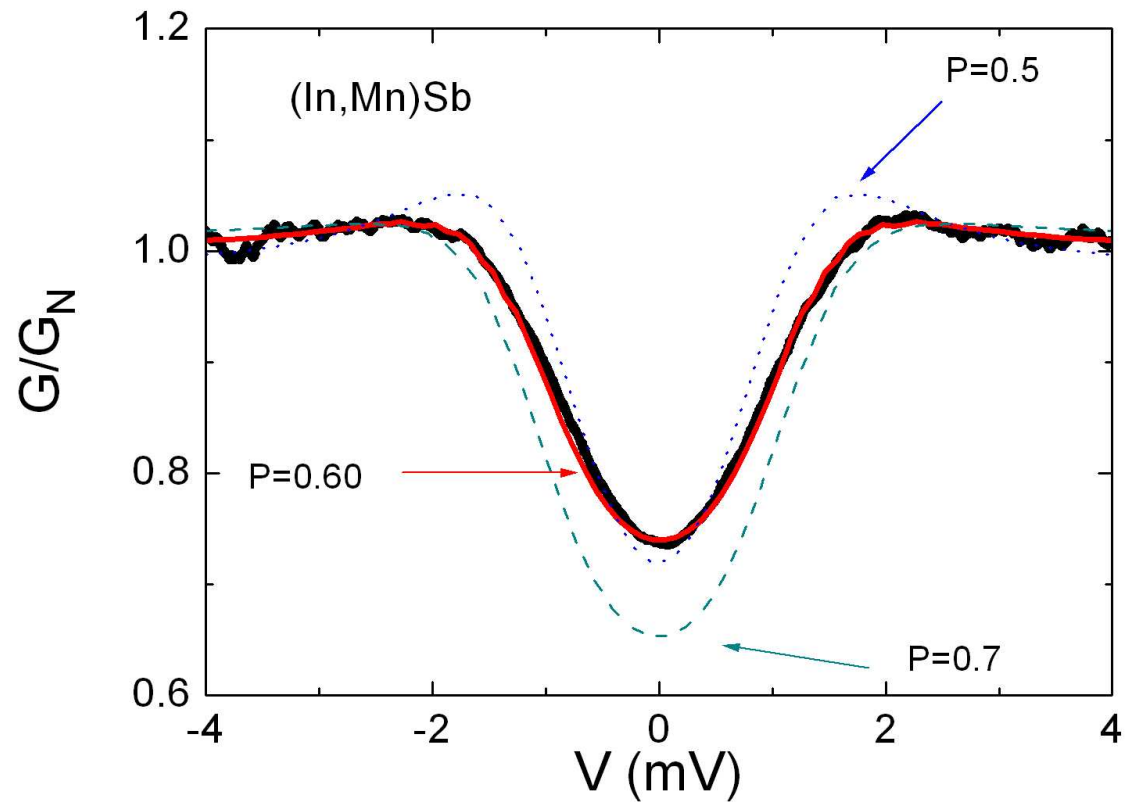
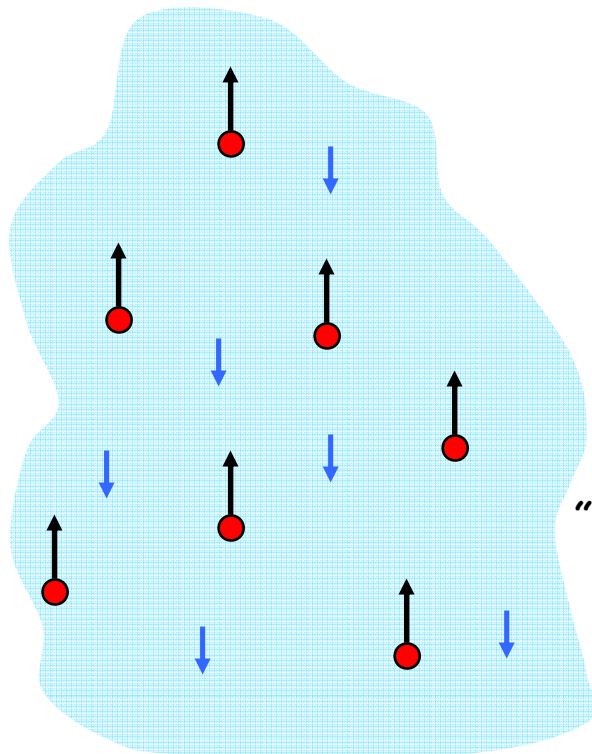
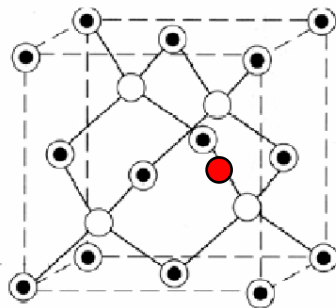
Néhány nanométeres kontaktus



Halbritter András
Geresdi Attila



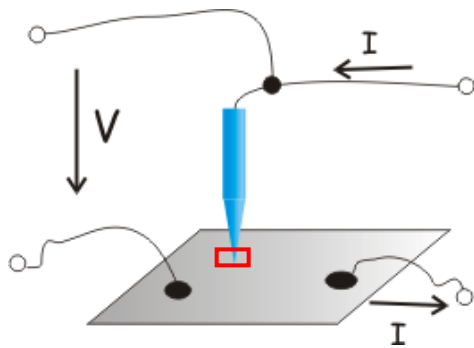
Spin-polarizáció (In,Mn)Sb mágneses félvezetőben



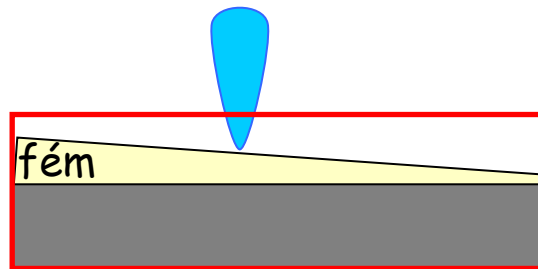
„carrier mediated magnetic coupling”

- M. Csontos *et al.*, Nature Materials **4**, 447 (2005)
- M. Csontos *et al.*, Phys. Rev. Lett. **95**, 227203 (2005)
- A. Geresdi *et al.*, Nano Letters (2008)
- G. Mihály *et al.*, Phys. Rev. Lett, March (2008)

Mágneses mikroszkópia

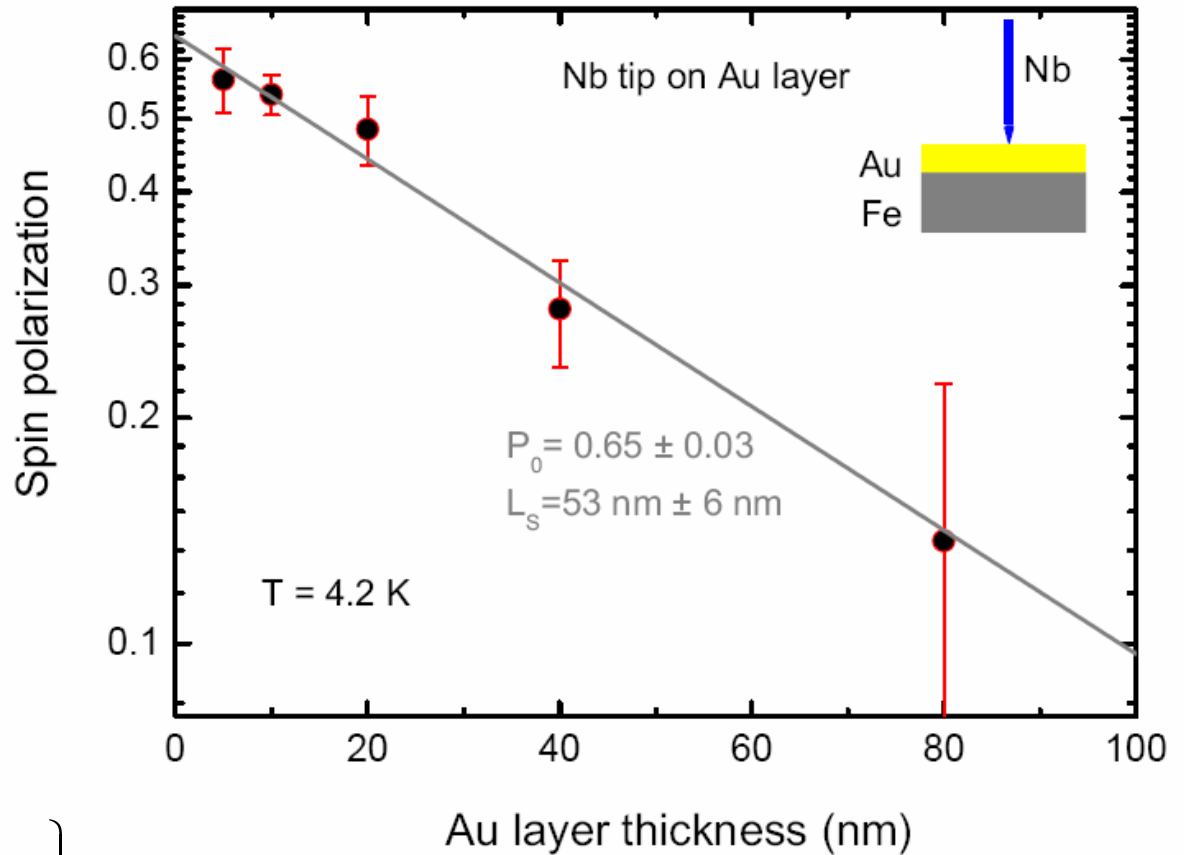


Spin diffúziós hossz

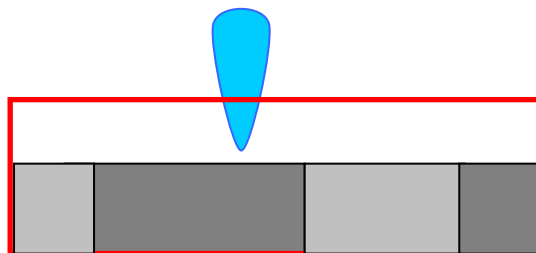


Mágneses fém

$$P = P_o \exp\left\{-\frac{l}{\sqrt{D\tau_s}}\right\}$$



Geresdi Attila, BME



Mágneses szerkezetek meghatározása
nanométeres méretskálán

Konklúzió

A klasszikus gondolatment alapján megértettnek hitt jelenség (pl. Hall effektus) igazi magyarázata már kvantumfizikai szemléletet feltételez.

Az alacsony hőmérsékleteken tapasztalt meglepő jelenségek (szupravezetés, kvantum Hall effektus) megértése sokszor egyszerűbb, mint a megszokott tulajdonságok korrekt leírása.

Az újonnan megfigyelt spin-függő jelenségek (pl. az anomális Hall-effektus mágneses félvezetőkben), igen gyorsan a spintronikai eszközök működési elvévé válnak.

Az Andrejev kísérletek felfedik az áram „mágnességét”. Ez a mennyiség meghatározó a nanoméretű mágneses memóriák írásánál – spin-forgatónyomaték.

<http://dept.phy.bme.hu/egzotikus.pdf>

Munkatársak



Kézmárki
István



Halbritter
András



Csonka
Szabolcs



Csontos
Miklós



Demkó
László



Bordács
Sándor



Makk
Péter



Geresdi
Attila

Ideális vezetés

Schrödinger egyenlet

$$-\frac{\hbar^2}{2m} \nabla^2 \psi = \varepsilon \psi$$

síkhullám

$$\psi(\mathbf{r}) = \frac{1}{\sqrt{V}} e^{i\mathbf{k}\mathbf{r}}$$

energia

$$\varepsilon(\mathbf{k}) = \frac{\hbar^2 \mathbf{k}^2}{2m}$$

impulzus

$$\mathbf{p} = \hbar \mathbf{k}$$

Áram

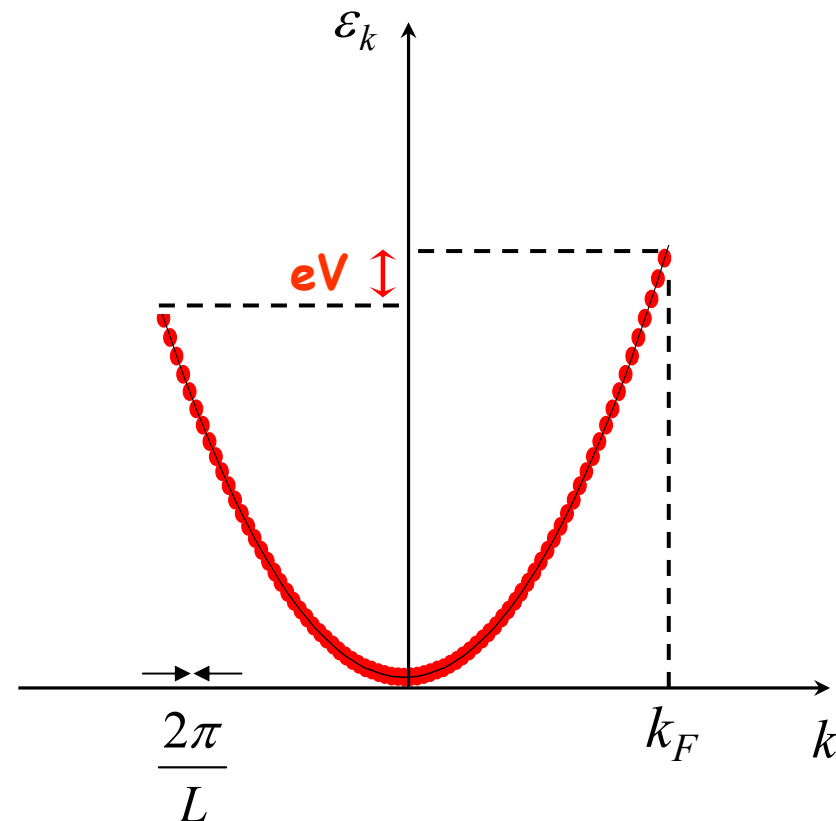
$$\mathbf{j} = ne\mathbf{v}$$

Sebesség

$$\mathbf{v} = \frac{\mathbf{p}}{m} = \frac{\hbar \mathbf{k}}{m} = \frac{1}{\hbar} \frac{\partial \varepsilon}{\partial \mathbf{k}}$$

Elektronok
sűrűsége

$$n = 2 \left(\frac{1}{L} \right) \frac{\partial \varepsilon}{\partial \mathbf{k}} \left(\frac{L}{2\pi} \right)$$



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Elektronok
sűrűsége

$$n = 2 \left(\frac{1}{L} \right) \frac{eV}{\frac{\partial \varepsilon}{\partial \mathbf{k}}} \left(\frac{L}{2\pi} \right)$$

$$\mathbf{j} = \frac{2e^2}{h} V \rightarrow G_0 = \frac{1}{12.9 \text{ k}\Omega}$$

Ideális vezetés

Schrödinger egyenlet

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síkhullám

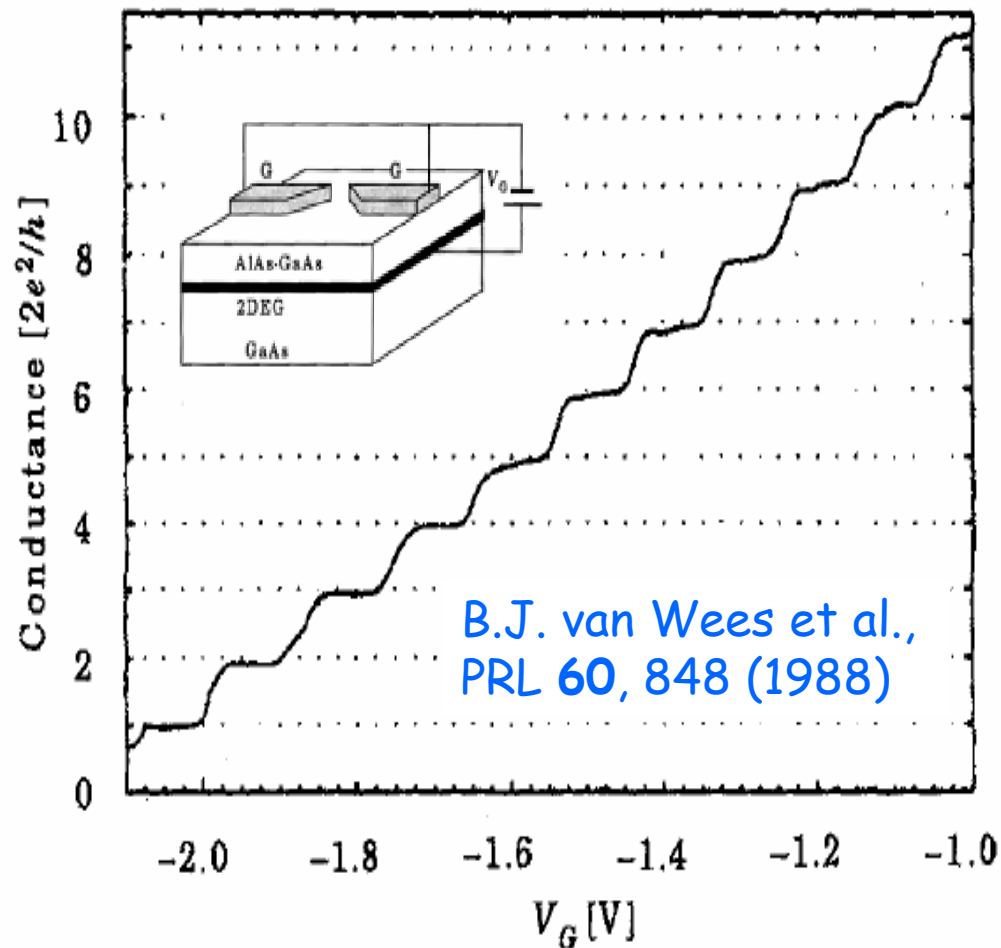
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A vezetési csatornák száma lépésekben változik a keresztmetszet növelésével

